
CORRECTIONS TO CCW PREDICTIONS FOR THE MOTION OF CYLINDRICAL HYDROMAGNETIC SHOCK WAVES

¹Dr Chandravir Singh, ²Dr R.S. Indolia & ³Dr Neera Sharma

¹²³Department of Physics, Agra College, Agra-28200

Email-drcvsinghph@gmail.com

ABSTRACT:

The Study of cylindrical hydro magnetic shock waves in presence of entropy effects have been study by CCW¹⁻³ predictions to improve their accuracy for the motion of diverging cylindrical hydro magnetic shock waves through in an ideal gas under its own entropy effect in the presence of constant axial magnetic field in presence of entropy .Assuming an initial density distribution $\rho_0 = \rho' r^{-w}$, where ρ' , is the density at the axis of symmetry and w , is constant, the analytical expression for flow variables representing both the situations via; weak and strong cases of shock have been obtained .

1. INTRODUCTION: Study of cylindrical hydro-magnetic shock waves has proved its importance in different branches of Science and Technology. In recent past Kumar⁴⁻⁶ ; and Kumar and Prakash⁷⁻⁸ have assuming initial density distribution $\rho_0 = \rho' r^{\pm w}$ have investigated the propagation of hydro magnetic cylindrical shock waves through axial magnetic field for both cases of weak and strong shock using CCW method . J.B. Singh and S.K.Pandey⁹ have studied the propagation of magnetogasdynamics cylindrical shock waves in self gravitating and rotating gas on the presence of constant axial magnetic field for both cases of weak and strong shock using CCW method. Assumption of any arbitrary initial density distribution of the form $\rho_0 = \rho' r^{\pm w}$ or $\rho_0 = \rho' \exp(\pm \lambda r)$ impose restriction values of propagation distance r , which are permitted by the fulfillment of initial entropy distribution condition i.e. $p_0 \rho_0^{-\gamma} = c^*$. In this paper ,the EOD on motion of diverging cylindrical hydro magnetic shock waves through an ideal gas under its own gravitation and rotation in presence of constant axial magnetic field simultaneously , for both (i) weak and (ii) strong cases of shock ,assuming an initial density decrease of the unperturbed medium as $\rho_0 = \rho' r^{-w}$.

2. BASIC EQUATIONS: The equations governing the flow of gas enclosed by the shock front are

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial r} + \frac{1}{\rho} \frac{\partial P}{\partial r} + \frac{\mu}{2\rho} \frac{\partial H^2}{\partial r} + \frac{Gm}{r^2} - \frac{v^2}{r} &= 0 \\ \frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial r} + \rho \left(\frac{\partial u}{\partial r} + \frac{u}{r} \right) &= 0 \\ \frac{\partial P}{\partial t} + u \frac{\partial P}{\partial r} - a^2 \left(\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial r} \right) &= 0 \quad (1) \\ \frac{\partial m}{\partial r} - 2\pi r \rho &= 0 \\ \left(\frac{\partial}{\partial t} + u \frac{\partial}{\partial r} \right) (vr) &= 0 \\ \frac{\partial H}{\partial t} + u \frac{\partial H}{\partial r} + H \frac{\partial u}{\partial r} + H \frac{u}{r} &= 0 \end{aligned}$$

Where r is the radial coordinate and r, v are the radial and azimuthal components of the particle velocity. P, ρ, H, μ and m denote respectively, the pressure, the density, the axial magnetic field, the magnetic permeability of the gas and mass inside the cylinder of unit cross-section of length r

3. BOUNDARY CONDITIONS: The magneto-hydrodynamic shock condition can be written in terms of single parameter $\xi = \rho/\rho_0$ as

$$\rho = \rho_0 \xi, \quad H = H_0 \xi, \quad u = \frac{\xi - 1}{\xi} U$$

(2)

$$U^2 = \frac{2\xi}{(\gamma+1) - (\gamma-1)\xi} \left[a_0^2 + \frac{b_0^2}{2} \{ (2-\gamma)\xi + \gamma \} \right] \text{ and}$$

$$P = P_0 + \frac{2\rho_0(\xi - 1)}{(\gamma + 1) - (\gamma - 1)\xi} \left[a_0^2 + \frac{\gamma - 1}{4} b_0^2 (\xi - 1)^2 \right]$$

Where 0 stands for the states immediately ahead of the shock, U is the shock velocity. a_0 is the sound speed $\sqrt{\gamma P_0/\rho_0}$ and b_0 is the alfvén speed $\sqrt{\mu H_0^2/\rho_0}$.

STRONG SHOCK: For the strong shock, ρ/ρ_0 is large. Now consider the two cases of weak and strong magnetic field.

Case I: For weak magnetic field $b_0^2 \ll a_0^2$, using this condition the boundary conditions (2) become:

$$\rho = \rho_0 \xi, \quad H = H_0 \xi, \quad u = \frac{\xi - 1}{\xi} U$$

$$\frac{P}{P_0} = 1 + \{ \chi' a_0^2 + A' b_0^2 \} \frac{U^2}{a_0^4}, \text{ where}$$

(3)

$$\chi' = \gamma(\xi - 1)/\xi, \quad A' = \frac{\gamma(\xi - 1)}{4\xi} [(\gamma - 1)(\xi - 1)^2 - 2\{(2 - \gamma)\xi + \gamma\}]$$

Case II: For strong magnetic field at $b_0^2 \gg a_0^2$ using this condition, the boundary condition (2) becomes

$$\rho = \rho_0 \xi, \quad H = H_0 \xi, \quad u = \frac{\xi - 1}{\xi} U$$

$$\frac{P}{P_0} = 1 + \chi \{ b_0^2 + A a_0^2 \} \frac{U^2}{a_0^2 b_0^2}, \text{ where}$$

(4)

$$\chi = \frac{\gamma(\gamma - 1)(\xi - 1)^3}{2\xi(2 - \gamma)(\xi + \gamma)}, \quad A = \frac{4}{(\gamma - 1)(\xi - 1)^2} - \frac{2}{(2 - \gamma)\xi + \gamma}$$

4. CHARACTERISTIC EQUATIONS: The characteristic form of the system of equation (1) is easily obtained by forming a linear combination of first and third equation of system of equation (1) an only one direction in (r, t) plane can be written as

$$dP + \mu H dH + \rho c du + \frac{\rho c^2 u dr}{(u+c)r} + \frac{\rho c G m}{(u+c)} \frac{dr}{r^2} - \frac{\rho c v^2}{(u+c)} \frac{dr}{r} = 0$$

(5)

In order to estimate the strength of overtaking disturbance an independent characteristic is considered .The differential relation valid across C_+ disturbance is written by replacing c by $-c$ in equation (5) we get

$$dP + \mu H dH - \rho c du + \frac{\rho c^2 u dr}{(u-c)r} + \frac{\rho c G m}{(u-c)} \frac{dr}{r^2} + \frac{\rho c v^2}{(u-c)} \frac{dr}{r} = 0$$

(6)

Where $c^2 = a^2 + b^2 = \frac{\gamma P}{\rho} + \frac{\mu H^2}{\rho}$.

5. ANALYTICAL RELATION FOR FLOW VARIABLES:

The equilibrium state of the gas is assumed to be specified by the condition $\partial/\partial t = 0$ and $H z_0 = \text{constant}$.

(7)

Using (7) the first equation of the system of the equation (1) the equilibrium prevailing condition in front of the shock can be written as

$$\frac{1}{\rho_0} \frac{dP_0}{dr} + \frac{Gm}{r^2} - \frac{V^2}{r} = 0$$

(8)

Assuming an initial density distribution law as $\rho_0 = \rho' r^{-w}$,where ρ' is the density at the axis of symmetry and w is constant.

Using $\rho_0 = \rho' r^{-w}$, from the fourth equation of the system of equation (1) we get

$$m = 2\pi \rho' r^{2-w} / (2-w) \quad (9)$$

From equation (7) and (8) we get

$$\frac{P_0}{P_1} = K + K_1 r^{2-w} + K_2 r^{1-2w}$$

$$\frac{da_0}{a_0} = \frac{1}{2} \left(\frac{dP_0}{P_0} + w \frac{dr}{r} \right) \quad (11)$$

Where K' is constant of integration $\frac{K'}{P_1} = K$, $K_1 = \frac{\Omega_0^2 \rho' }{(2-w)P'}$, $K_2 = \frac{2\pi\gamma}{D(1-2w)(w-2)}$, $v = \Omega_0 r \text{ and}$

$D = a'^2 / G \rho' P'$ is the pressure at the axis and a' is the sound velocity at the axis .

5(a). STRONG SHOCK WITH WEAK MAGNETIC FIELD (SSWMF):

By substituting the shock conditions (3) in to equation (5) and respective values of various quantities, we get

$$\frac{dU^2}{dr} + U^2 \left[\frac{M_1}{r} + M_2 \beta^2 r^{1-w} + M_3 \beta^2 r^{-2w} + M_4 \beta^2 r^{3-2w} + M_5 \beta^2 r^{2-3w} + M_6 \beta^2 r^{1-4w} \right] + M_7 r^{-w} + M_8 r + M_9 \beta^2 r^{1-3w} - M_{10} r^{3-3w} + M_{11} \beta^2 r^{2-2w} = 0 \quad (12)$$

Where $M = \left[\frac{\chi'}{\gamma} + \frac{\xi'}{2} \sqrt{\frac{\chi'}{\xi}} \right]$, $N = \left[\frac{(\xi-1)}{(\xi-1) + \sqrt{\chi'\xi}} - \frac{w}{\gamma} \right] \frac{\chi'}{m}$, $N_1 = \left[\frac{w}{\gamma} + \frac{N}{M\gamma} \right] \frac{A'\beta^2}{MK}$, $M_1 = N - N_1$

$$M_2 = \left[\frac{N}{M} - 2(1-w) \right] \frac{A'K_1}{M\gamma K^2} \quad , \quad M_3 = \frac{A'K_2}{M\gamma K^2} \left[(3w-1) + \frac{N}{M} \right]$$

$$M_4 = 2A'K_1^2 (2-w) / M\gamma K^3 \quad , \quad M_5 = 6K_1 K_2 (1-w) / M\gamma K^3$$

$$M_6 = 2K_2^2 (1-2w) A' / M\gamma K^3 \quad , \quad M_7 = N_2 - N_3 \quad , \quad M_8 = N_5 - N_4 \quad ,$$

$$M_9 = N_2 K_2 A' / M\gamma K^2, \quad M_{10} = N_4 A' / M\gamma K^2, \\ M_{11} = A'(N_2 K_1 - N_4 K_2) / M\gamma K^2 \\ N_2 = \frac{2\pi a^2}{D(2-w)M} \left[\frac{\xi\sqrt{\chi\xi}}{\{(\xi-1)+\sqrt{\chi\xi}\}} - 1 \right], \quad N_3 = \beta^2 A' N_2 / M\gamma K, \\ N_4 = \frac{\Omega^2}{M} \left[\frac{\xi\sqrt{\chi\xi}}{\{(\xi-1)+\sqrt{\chi\xi}\}} - 1 \right], \quad N_5 = \beta^2 A' N_4 / M\gamma K,$$

On integration equation (27) we get

$$\frac{U^2}{F_P^{2w}} = [K_p r^{-M_1} - M_{12} r^{1-w} - M_{13} r^{-2} - M_{14} \beta^2 r^{2-3w} + M_{15} \beta^2 r^{4-3w} - M_{16} \beta^2 r^{3-2w} - M_{17} \beta^2 r^{5-3w} - M_{18} \beta^2 r^{4-4w} - M_{19} \beta^2 r^{3-5w} - M_{20} r^{4-w} + M_{21} \beta^2 r^6 - 2w = 0] \quad (13)$$

Exp. { $\beta^2 (M_{22} r^{2-w} + M_{23} r^{1-2w} + M_{24} r^{4-2w} + M_{25} r^{3-3w} + M_{26} r^{2-4w})$ } where K_p is constant of integration and $M_{12} \dots M_{26}$ are other constants.

For the C. disturbance generated by the shock, the fluid velocity increment using may be expressed as

$$du = \frac{\xi-1}{2\xi} \left[U \left(\frac{M_1}{r} + M_2 \beta^2 r^{1-w} + M_3 \beta^2 r^{-2w} + M_4 \beta^2 r^{3-2w} + M_5 \beta^2 r^{2-3w} + M_6 \beta^2 r^{1-4w} \right) + \frac{1}{U} (M_7 r^{-w} + M_8 r + M_9 \beta^2 r^{1-3w} - M_{10} r^{3-3w} + M_{11} \beta^2 r^{2-2w}) \right] dr \quad (14)$$

Substituting the shock condition (3) in to equation (6) and using the shock condition (3), we get

$$du_+ = -\frac{\xi-1}{2\xi} \left[U \left(\frac{H_1}{r} + H_2 \beta^2 r^{1-w} + H_3 \beta^2 r^{-2w} + H_4 \beta^2 r^{3-2w} + H_5 \beta^2 r^{2-3w} + H_6 \beta^2 r^{1-4w} \right) + \frac{1}{U} (H_7 r^{-w} + H_8 r + H_9 \beta^2 r^{1-3w} - H_{10} r^{3-3w} + H_{11} \beta^2 r^{2-2w}) \right] dr \quad (15)$$

Now in presence of both C_+ and C. disturbances the fluid velocity increment behind the shock will be related as

$$du + du_+ = \frac{\xi-1}{\xi} dU \quad (16)$$

using equation (14),(15) and relation(16) we get

$$\frac{dU^2}{dr} + U^2 \left[\frac{S_1}{r} + S_2 \beta^2 r^{1-w} + S_3 \beta^2 r^{-2w} + S_4 \beta^2 r^{3-2w} + S_5 \beta^2 r^{2-3w} + S_6 \beta^2 r^{1-4w} \right] + S_7 r^{-w} + S_8 r + S_9 \beta^2 r^{1-3w} - S_{10} r^{3-3w} + S_{11} \beta^2 r^{2-2w} = 0 \quad (17)$$

Where $S_1 = H_1 + M_1, S_2 = M_2 + H_2, S_3 = M_3 + H_3, S_4 = M_4 + H_4, S_5 = M_5 + H_5, S_6 = M_6 + H_6, S_7 = M_7 + H_7, S_8 = M_8 + H_8, S_9 = M_9 + H_9, S_{10} = M_{10} + H_{10}, S_{11} = M_{11} + H_{11}$ on integration, we get

$$U^2_{EOD} = [K_p^* r^{-S_1} - S_{12} r^{1-w} - S_{13} r^{-2} - S_{14} \beta^2 r^{2-3w} + S_{15} \beta^2 r^{4-3w} - S_{16} \beta^2 r^{3-2w} - S_{17} \beta^2 r^{5-3w} - S_{18} \beta^2 r^{4-4w} - S_{19} \beta^2 r^{3-5w} - S_{20} r^{4-w} + S_{21} \beta^2 r^6 - 2w] \exp. \{ \beta^2 (S_{22} r^{2-w} + S_{23} r^{1-2w} + S_{24} r^{4-2w} + S_{25} r^{3-3w} + S_{26} r^{2-4w}) \}$$

where K_p^* is constant of integration and $S_{12} \dots S_{26}$ are other constants.

(18)

5(b). STRONG SHOCK WITH STRONG MAGNETIC FIELD

(SSSMF):

By substituting the shock conditions (4) in to equation (5) and respective values of various quantities, we get

$$\frac{dU^2}{dr} + U^2 \left[\frac{B_1}{r} + B_2 \beta^{-2} r^{-2w} + B_3 \beta^{-2} r^{1-w} \right] + B_4 r^{-w} + B_5 r - B_6 \beta^{-2} r^{1-3w} + B_7 \beta^{-2} r^{3-w} + B_8 \beta^{-2} r^{2-2w} = 0 \quad (19)$$

Where $B = \left[\frac{\chi}{\gamma} + \frac{\xi-1}{2} \sqrt{\frac{\chi}{\xi}} \right], Z = \left[\frac{\chi(\xi-1)}{B\{(\xi-1)+\sqrt{\chi\xi}\}} - \frac{\chi w}{B\gamma} \right], Z_1 = \left[\frac{\chi w}{\gamma B} - Z \right] \frac{\chi A \beta^2 K}{B\gamma}, B_1 = Z + Z_1$

$$B_2 = \frac{A\chi K_2}{B\gamma} \left[(1 - 3w) + \frac{w\chi}{\gamma B} - \frac{\chi(\xi-1)}{B\{(\xi-1)+\sqrt{\chi\xi}\}} + \frac{\chi w}{B\gamma} \right], Z_5 = \frac{Z_4 \chi A \beta^2 K}{B\gamma}$$

$$B_3 = \frac{A\chi K_1}{B\gamma} \left[(2 - 2w) + \frac{w\chi}{\gamma B} - \frac{\chi(\xi-1)}{B\{(\xi-1)+\sqrt{\chi\xi}\}} + \frac{\chi w}{B\gamma} \right], Z_3 = \frac{Z_2 \chi A \beta^2 K}{B\gamma}$$

$$B_4=Z_2-Z_3, B_5=Z_5-Z_4, B_6=\frac{Z_2 \chi^A K_2}{B_\gamma}, B_7=\frac{Z_4 \chi^A K_1}{B_\gamma}, B_8=\frac{\chi^A [Z_4 K_2 - Z_2 K_1]}{B_\gamma}$$

$$Z_2=\frac{2\pi a^2}{D(2-w)B} \left[\frac{\xi \sqrt{\chi \xi}}{\{(\xi-1)+\sqrt{\chi \xi}\}} - 1 \right], Z_4=\frac{\Omega^2}{B} \left[\frac{\xi \sqrt{\chi \xi}}{\{(\xi-1)+\sqrt{\chi \xi}\}} - 1 \right]$$

On integration equation(19) ,we get

$$\frac{U^2}{F_P} = [K_s r^{-B_1} - B_9 r^{1-w} - B_{10} r^2 - B_{11} \beta^{-2} r^{2-3w} - B_{12} \beta^{-2} r^{4-2w} - B_{13} \beta^{-2} r^{3-2w}] \exp. \{ \beta^2 (B_{14} r^{1-2w} + B_{15} r^{2-w}) \} (20)$$

where K_s is constant of integration and B_9, \dots, B_{15} are other constants.

For the C. disturbances generated by the shock ,the fluid velocity increment using may be expressed as

$$du_- = -\frac{(\xi-1)}{2\xi} \left[\left(\frac{B_1}{r} + B_2 \beta^{-2} r^{-2w} + B_3 \beta^{-2} r^{1-w} \right) U + \frac{1}{U} + (B_4 r^{-w} + B_5 r - B_6 \beta^{-2} r^{1-3w} + B_7 \beta^{-2} r^{3-w} + B_8 \beta^{-2} r^{2-2w}) \right] dr$$

(21)

Substituted the shock condition (3) in to equation(6) and again using shock condition(3), we get

$$du_+ = \frac{(\xi-1)}{2\xi} \left[\left(\frac{B_1}{r} + B_2 \beta^{-2} r^{-2w} + B_3 \beta^{-2} r^{1-w} \right) U + \frac{1}{U} + (B_4 r^{-w} + B_5 r - B_6 \beta^{-2} r^{1-3w} + B_7 \beta^{-2} r^{3-w} + B_8 \beta^{-2} r^{2-2w}) \right] dr$$

(22)

Now in presence of both C_+ and C. disturbances the fluid velocity increment behind the shock will be related as

$$du_- + du_+ = \frac{\xi-1}{\xi} dU \tag{23}$$

using equation(21),(22) and (23) we get

$$\frac{dU^2}{dr} + \left[\frac{h_1}{r} + h_2 \beta^{-2} r^{-2w} + h_3 \beta^{-2} r^{1-w} \right] U^2 + h_4 r^{-w} + h_5 r - h_6 \beta^{-2} r^{1-3w} + h_7 \beta^{-2} r^{3-w} + h_8 \beta^{-2} r^{2-2w} = 0 \tag{24}$$

Where $h_1 = B_1 + L_1, h_2 = B_2 + L_2, h_3 = B_3 + L_3, h_4 = B_4 + L_4, h_5 = B_5 + L_5, h_6 = B_6 + L_6, h_7 = B_7 + L_7, h_8 = B_8 + L_8$

On integration (24) we get

$$U^2_{EOD} = [K_s^* r^{-h_1} - h_9 r^{1-w} - h_{10} r^2 - h_{11} \beta^{-2} r^{2-3w} + h_{12} \beta^{-2} r^{4-2w} - h_{13} \beta^{-2} r^{3-2w}] \exp. \{ \beta^{-2} (h_{14} r^{1-2w} + h_{15} r^{2-w}) \} \tag{25}$$

where K_s^* is constant of integration and

$$h_9 = h_4 / (1 + h_1 - w), \quad h_{10} = h_5 / (2 + h_1), \quad h_{11} = [h_2 h_4 / (1 - 2w) - h_6] / (2 + h_1 - 3w), \quad h_{12} = [h_3 h_5 / (2 - w) + h_7] / (4 + h_1 - w)$$

$$h_{13} = [h_2 h_5 / (1 - 2w) + h_3 h_4 / (2 - w) + h_8] / (3 + h_1 - 2w), \quad h_{14} = h_2 / (2w - 1), \quad h_{15} = h_3 / (w - 2)$$

The flow variable expressions F.P. and EOD can be written by substituting (13,20) and (18,25) respectively in (3) and (4) and use for computation.

Permissible Shock Front Location (pslf): The pressure and the density in unperturbed state given by expression (13) and $\rho_o = \rho' r^{-w}$ must fulfill the initial entropy distribution condition i.e. $P_o \rho_o^\gamma - c^*$ where C^* is constant, substituting the respective expression for P_o and ρ_o , we get

$$K_1 r^{w\gamma} + K_2 r^{2-w(1-\gamma)} + K_3 r^{1-w(1-\gamma)} - \frac{c^* \rho'^\gamma}{p'} = 0 \tag{26}$$

This is an equation in different powers depending upon the value of w and γ of propagation distance r it may have number of roots equal to highest powers of r , in general. The roots have been evaluated by expression (26). Hit and trial method.

RESULT AND DISCUSSION: Numerical estimates of flow variables for both F.P. and having included the EOD have been computed only at those locations of the shock front which are permitted by the initial entropy distribution condition in the unperturbed state. The results have also shown the comparison with F.P. predictions (a)(i) taking $U/a_0=4.098543, 4.38539, 4.583945, 4.78935, 5.01525$ at $r=0.45, \beta^2=0.10, 0.15, \Omega^2=10, 15$ $w=1.0, 1.05, \gamma=1.4, D=0.2, 0.3, 0.4, \xi=1.5$ (ii) taking $U/a_0=12.5436$ at $r=0.55, \beta^2=1.5, \gamma=1.4, w=1.0, D=0.2, \xi=3$ (iii) taking $U/a_0=20.0568$ at $r=0.55, \beta^2=1.5, \gamma=1.4, w=1.0, D=0.2, \xi=4.5$ (iv) taking $U/a_0=120.3856$ at $r=0.55, \beta^2=1.5, \gamma=1.4, w=1.0, D=0.2, \xi=5.9$ SSWMF (b)(i) taking $U/a_0=7.80538, 9.5869, 9.58685, 9.78345, 9.98523$ at $r=0.55, \beta^2=10, 15, \Omega^2=10, 15, w=1.0, 1.05, \gamma=1.4, D=0.2, 0.3, 0.4, \xi=1.5$ (ii) taking $U/a_0=17.54369$ at $r=0.55, \beta^2=15, \gamma=1.4, w=1.0, D=0.2, \xi=3$ (iii) taking $U/a_0=5.89546$ at $r=0.55, \beta^2=15, \gamma=1.4, w=1.0, D=0.2, \xi=4.5$ (iv) taking $U/a_0=140.58395$ at $r=0.55, \beta^2=15, \gamma=1.4, w=1.0, D=0.2, \xi=5.9$ SSSWF, numerical estimates of flow variables have been computed at permissible propagation distance for both F.P. and having included the EOD. It is observed from comparison those numerical values of flow variables, representing free propagation with the present values of flow variables included the EOD that the EOD over all retains the qualitative variation with parameter $r, \beta^2, \Omega^2, \xi, w$ and D unchanged.

REFERENCES

1. Chisnell, R.F. :Proc. R.Soc. London, 232 A, 350 (1955).
2. Chester, W: Phil. Mag. 45, 7, 1293 (1954).
3. Whitham, G.B.: J. Fluid, Mech., 4, 337 (1958).
4. Kumar. S and Singh C.V.: J. Ultra. Sci. Phy. Sci. 16, (1), 77, 2004.
5. Kumar. S and Saxena, D.K.: Astrophys. space sci. 100, 65 (1984).
6. Kumar. S and Kishor B.: J.P.A.S. 6, 156 (2000).
7. Kumar. S and Prakash R.: II. Nuovo Ciminto 77, B, 191 (1983).
8. Kumar. S and Prakash R : Astrophys. space sci. 235 27 (1996).
9. Singh JB and Pandey, S.K.: Astrophys. space sci. 141 221 (1988).
10. Yousaf, M : J. Fluid, Mech., 66 577 (1974).
11. Yousaf, M : Phys. Fluid, 25 (1), 45, (1982).
12. Yousaf, M : Phys. Fluid, 28 1659 (1985).